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**MIXING RULES FOR AND EFFECTS OF OTHER HYDROGEN ISOTOPES AND OF
ISOTOPIC SWAPPING ON TRITIUM RECOVERY AND LOSS TO BIOSPHERE FROM FUSION
REACTORS**

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ABSTRACT

Efficient recovery of bred and unburnt tritium from fusion reactors, and control of its migration within reactors and of its escape into the biosphere are essential for self-sufficient fuel cycles and for public, plant personnel, and environmental protection. Tritium in fusion reactors will be mixed with other isotopes and protium, interacting to form reactions and diffusion into coolant loops from steam cycles. Rational design for tritium recovery and escape prevention must acknowledge this fact. Consequences of isotopic admixture are explored, mixing rules for proposed fusion reactor dilute-solution conditions are developed, and a table of tritium recovery factors is presented.

INTRODUCTION

Present indications are that at least the first generation of fusion reactors will be filled with a mixture of deuterium and tritium (D-T), because it is considered that such a mixture will induce sufficient neutrons to maintain the reaction, less extensive than that required for other processes. The D-T mixture is being developed, however, to be used in the gas phase, and the breeding of tritium from deuterium in these reactors is expected. The breeding of tritium from deuterium in the reaction $D + T \rightarrow He + n$ is expected to be 50%, though tritium is also bred in other reactions, relatively minor, such as $D + D \rightarrow He + T + n$ and $D + D \rightarrow He + T + \gamma$, and $D + D \rightarrow He + T + n$, and $D + D \rightarrow He + T + n$. Logical breeding of tritium is expected.

Efficient recovery of bred and unburnt tritium from fusion reactors, and control of its migration within reactors and of its escape into the biosphere are essential for self-sufficient fuel cycles and for public, plant personnel, and environmental protection.

Work done under the auspices of the U.S. AEC.

plants is essential for management of volatile fission product tritium, and for preventing the escape of natural sources of tritium to the biosphere. It is expected that tritium will be bred in fusion reactors, and that it will be used in fusion reactors.

Present indications are that at least the first generation of fusion reactors will be filled with a mixture of deuterium and tritium (D-T), because it is considered that such a mixture will induce sufficient neutrons to maintain the reaction, less extensive than that required for other processes. The D-T mixture is being developed, however, to be used in the gas phase, and the breeding of tritium from deuterium in these reactors is expected. The breeding of tritium from deuterium in the reaction $D + T \rightarrow He + n$ is expected to be 50%, though tritium is also bred in other reactions, relatively minor, such as $D + D \rightarrow He + T + n$ and $D + D \rightarrow He + T + \gamma$, and $D + D \rightarrow He + T + n$, and $D + D \rightarrow He + T + n$. Logical breeding of tritium is expected. Efficient recovery of bred and unburnt tritium from fusion reactors, and control of its migration within reactors and of its escape into the biosphere are essential for self-sufficient fuel cycles and for public, plant personnel, and environmental protection.

Efficient recovery of bred and unburnt tritium is essential for maintaining low tritium inventories. Escape of tritium to the biosphere can be minimized by maintaining low tritium inventories and by providing engineered barriers to tritium escape. Under emergency conditions, the breeding blanket and the main coolant loops are usually considered to represent the greatest hazard, whereas during normal operations, thin, hot reactor components lying across direct paths to the biosphere are critical to tritium escape. Tritium escape into steam cycle loops, and thence into the biosphere, is generally regarded to be the most important pathway to be kept under control with normal operation. Operations in power plants employing a steam cycle, and the associated tritium recovery systems and escape-prevention barriers.

Protium will be generated by fission reactions of ^{235}U and ^{239}Pu in the reactor core. The tritium inventory in the reactor core will be maintained at a level which is sufficient to meet the needs of the breeding blanket and the main coolant loops. The tritium inventory in the breeding blanket and the main coolant loops will be maintained at a level which is sufficient to meet the needs of the breeding blanket and the main coolant loops.

Both the breeding blanket and the main coolant loops will be equipped with tritium recovery systems. The breeding blanket will be equipped with a tritium recovery system which will recover tritium from the breeding blanket and the main coolant loops. The main coolant loops will be equipped with a tritium recovery system which will recover tritium from the breeding blanket and the main coolant loops.

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and deuterium. Thus, unburnt tritium and deuterium, which constitute the majority of fuel fed to fusion reactors, may not represent the majority of hydrogen isotopes recovered.

Recovery of deuterium, which is relatively abundant in nature and can be obtained at modest cost, is not necessary for economic generation of fusion energy. However, buildup of protium and deuterium within reactor subsystems to excessive levels cannot be tolerated, because of potential deleterious effects on reactor materials of construction. Tritium recovery systems will simultaneously remove protium and deuterium because of their similar physicochemical characteristics, but at a penalty of increased tritium recovery system capacity.

Protium concentrations in thermonuclear reactor fuel streams will probably have to be reduced to 1% or less. The D/T isotopic ratio must be adjusted by the addition of medium fractionation.

Clearly, external design procedures for fusion plant systems that use part of the tritium fuel cycle are affected by the presence of tritium, and resulting assessment of its hazards must take into account the effects of the presence of the other fuel isotopes. The tritium inventory in the breeding blanket and the main coolant loops will be maintained at a level which is sufficient to meet the needs of the breeding blanket and the main coolant loops.

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Our discussion will be confined to the dilute-solution conditions anticipated for tritium in fusion reactor breeding blankets, coolant loops, structural materials, and perhaps cavity exhaust streams. Early estimates of rates of tritium escape into steam cycles from coolant loops suggested that tritium concentrations in fusion reactor systems might have to be maintained at 1 ppm or less. As the difficulties of tritium recovery from solutions at such concentrations have become apparent, more effective concepts for prevention of tritium escape have been developed, and the importance of potential tritium hazards relative to other potential fusion reactor hazards has been more accurately assessed, so that a justification has been given to operation with tritium concentrations an order or two higher in magnitude. Nonetheless, dilute solution theory will still be applicable, at least up to a higher concentration than one would expect, and still quite low.

We will focus on the development of mixing rules for equilibria of hydrogen isotope mixtures, and on the effects that the presence of more than one hydrogen isotope has on the properties of the fusion reactor systems.

MIXING RULES FOR THE EQUILIBRIA OF HYDROGEN ISOTOPE MIXTURES IN METALS

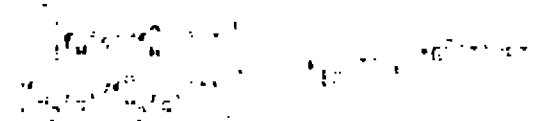
Most of the structure of fusion reactors will probably be metallic. Many of the reactions that involve liquid metals, such as the breeding of tritium from deuterium, and primary coolant loops. Clearly, the solubility of deuterium, as affected by the presence of other hydrogen isotopes, is critical to the design of tritium breeding and the performance of the tritium recovery systems and associated hardware. We will derive mixing rules for the equilibria of hydrogen isotope mixtures in metals.

Review of Equilibrium Thermodynamics of Solubility and Diffusion in Metals

Many careful studies have been made by Sieverts,⁽⁹⁾ Ricca,⁽¹⁰⁾ Pines,⁽¹¹⁾ and others. In addition, the work of Hank,⁽¹²⁾ and Wu,⁽¹³⁾ and Sieverts and Vanhook,⁽¹⁴⁾ have confirmed that the empirical

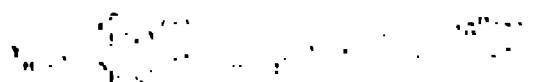
relationship named for A. Sieverts accurately describes the temperature-dependent solubility of protium, deuterium, and tritium in metals at very low concentrations and, in some instances, at relatively high concentrations, for both pure metals and alloys with low surface-to-volume ratios so that adsorption is not a significant factor.

A theoretical justification for Sieverts' empirical law can be developed as follows. We assume, in accord with experimental evidence, e.g., see Isenberg,⁽¹⁵⁾ that a hydrogen solution exists in solution in metals as an interstitial solution. For present purposes, it is not necessary whether the dissociation of dihydrogen is necessary for dissolution occurs in the gas phase, on the metallic surface, or in the bulk metal. An equilibrium expression for the interstitial solution is



where $H_{int}(M)$ is the interstitial hydrogen in metal.

The equilibrium constant for this reaction is $K = \frac{a_{H_{int}(M)}^2}{p_{H_2}}$, where $a_{H_{int}(M)}$ is the activity of interstitial hydrogen in metal, and p_{H_2} is the partial pressure of hydrogen gas. The activity of interstitial hydrogen in metal is defined as $a_{H_{int}(M)} = \gamma_{H_{int}(M)} x_{H_{int}(M)}$, where $\gamma_{H_{int}(M)}$ is the activity coefficient of interstitial hydrogen in metal, and $x_{H_{int}(M)}$ is the mole fraction of interstitial hydrogen in metal. The activity coefficient of interstitial hydrogen in metal is defined as $\gamma_{H_{int}(M)} = \frac{a_{H_{int}(M)}}{x_{H_{int}(M)}}$.



where K is the equilibrium constant for the reaction, $\gamma_{H_{int}(M)}$ is the activity coefficient of interstitial hydrogen in metal, $x_{H_{int}(M)}$ is the mole fraction of interstitial hydrogen in metal, and p_{H_2} is the partial pressure of hydrogen gas.

If temperature is constant, the law that the standard enthalpy of solution is constant for any metal

ed, then Stevarts' law for the solubility of hydrogen isotopes in metals is obtained:

$$X_{H(s)} = K_{S_H} (T) \left[P_{H_2} / 1 \text{ atm} \right]^{1/2}$$

$$K_{S_H} (T) = K_{S_H}^0 e^{-Q_{S_H}/RT} e^{-\frac{\Delta S_H^0}{R} \ln \frac{T}{T_0}}$$

$$Q_{S_H} = \Delta H^0$$

Derivation of Stevarts' law

It is well known that the only other explicit consideration of the problem of predicting solubility of the individual hydrogen isotopes in a mixture and the gross solubility of the mixture in a metal were suggested by James of Hickman [1] and by Stevarts [2]. James and Hickman suggested that gas atoms which do not should be considered as a single chemical species and that the solubility of a mixture of Stevarts' law is a simple linear combination:

$$X_{Total} = \sum_{i=1}^n X_{i(s)} = \sum_{i=1}^n K_{S_i} (T) P_i^{1/2}$$

He further suggested that the solubility of a mixture of isotopes is given by:

$$K_{S_{Total}} (T) = \sum_{i=1}^n K_{S_i} (T) = \sum_{i=1}^n K_{S_i}^0 e^{-Q_{S_i}/RT} e^{-\frac{\Delta S_i^0}{R} \ln \frac{T}{T_0}}$$

Hickman suggested that only a limited number of holes in the lattice are available for a mixture of isotopes and that the solubility of a mixture of isotopes is given by the sum of the solubilities of the individual isotopes in the lattice.

Hickman also presented experimental data for the solubility of a mixture of isotopes of protium and deuterium in a metal alloy of palladium and silver. On the basis of these experimental results, he concluded that the postulated mixing rule described the solubility of the mixture acceptably well.

Adkins and Gurevich used the mixing rule suggested by Hickman in analyzing the results of

experiments which investigated the permeation of gaseous mixtures through the walls of tubes of the same alloy. They claimed acceptable agreement between their experimental results and those of analysis based on Hickman's suggested mixing rule.

However, the fact that Stevarts' law accurately predicts the solubility of hydrogen isotopes in many metals and allows up to relatively high concentrations, suggests, for systems sufficiently dilute in the dissolved hydrogen isotopes, that the systems are also relatively dilute in comparison to the number of holes in the lattice valence bands. The large isotopic differences in solubility, as reflected by differences in Stevarts' law constants, suggest that a mixture of hydrogen isotopes cannot be treated as a single chemical substance for all purposes. Therefore, the following analysis, which explores the consequences of validity of Stevarts' law for dissolution of individual hydrogen isotopes in metals, was performed to develop an alternative, which has some theoretical basis, to Hickman's isotopic mixing rule.

On an atomic basis, in solutions sufficiently dilute in all solutes, solute atoms are surrounded by and interact with virtually all the time with solvent atoms rather than with other solute atoms. Therefore, in such solutions the solubility of a solute is more likely to be determined primarily by its properties and by those of the solvent and to be affected minimally by the properties of other solutes. This interpretation should also be valid if the solvent is a mixture of species, as long as the ratios of concentrations of the species in the solvent remain constant, e.g., if the solvent is an alloy, dilute or otherwise.

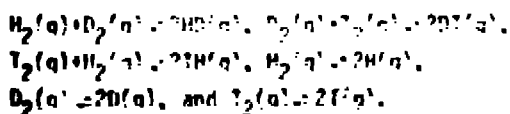
Thus we conclude that Stevarts' law expressions obtained by fitting experimental solubility data for pure isotopes are applicable to individual isotope solubilities when isotopic mixtures are involved. Therefore, for a system containing all three hydrogen isotopes, we write:

$$X_{Total} = X_{H(s)} + X_{D(s)} + X_{T(s)} \\ = K_{S_H} (T) P_{H_2}^{1/2} + K_{S_D} (T) P_{D_2}^{1/2} + K_{S_T} (T) P_{T_2}^{1/2}$$

$$= \left[K_{S_H}(T)^{1/2} K_{S_D}(T)^{1/2} K_{S_T}(T)^{1/2} \right] P_{Total}^{1/2}$$

$$= K_{S_{Total}}(T) P_{Total}^{1/2}$$

In an isotopic mixture there will be isotopic exchange in the gas phase, i.e., in addition to the molecular species $H_2(g)$, $D_2(g)$, and $T_2(g)$ the species $HD(g)$, $DT(g)$, and $TD(g)$ will be present. In addition, the diatomic molecules can dissociate to yield the atomic species $H(g)$, $D(g)$, and $T(g)$. The following set of six gas-phase reactions are independent:



The corresponding equilibria can be expressed in the forms:

$$\frac{P_{HD}^2(g)}{P_{H_2}(g) P_{D_2}(g)} = \frac{v_{HD}^2}{v_{H_2} v_{D_2}} = K_{HD}(T)$$

$$\frac{P_{HT}^2(g)}{P_{H_2}(g) P_{T_2}(g)} = \frac{v_{HT}^2}{v_{H_2} v_{T_2}} = K_{HT}(T)$$

with similar expressions for the other reactions. Experimental values for the equilibrium constants are listed in Table I.

In terms of y_H , y_D , and y_T the equilibrium relationships for the gas phase can be written, in the absence of significant molecular dissociation to form the atomic species, in the form:

$$\frac{\left(\frac{v_H}{v_{H_2}} - y_{H_2}^{1/2} + \frac{v_D}{v_{D_2}} - y_{D_2}^{1/2} + \frac{v_T}{v_{T_2}} - y_{T_2}^{1/2} \right)^2}{y_{H_2} y_{D_2}} = K_{HD}(T)$$

upon substitution of dissociation equilibrium constants as:

$$v_{HD} = \left(\frac{v_H}{v_{H_2}} - y_{H_2}^{1/2} + \frac{v_D}{v_{D_2}} - y_{D_2}^{1/2} + \frac{v_T}{v_{T_2}} - y_{T_2}^{1/2} \right) v_{H_2} y_{D_2}^{1/2}$$

Solution for the mole fractions of the atomic species may have to be obtained numerically. Substitution of the appropriate values for the diatomic species into the mixing rule then gives an insight into the relative values of v_H , v_D , and v_T .

If dissociation is not negligible, then the form of the atomic species equilibrium constants, in the situation of a gas mixture, is more complex.

TABLE I

EQUILIBRIUM CONSTANTS $K_{HD}(T)$, $K_{HT}(T)$, $K_{DT}(T)$, $K_{HD}(T)$, $K_{HT}(T)$, $K_{DT}(T)$, $K_{HD}(T)$, $K_{HT}(T)$, $K_{DT}(T)$, $K_{HD}(T)$, $K_{HT}(T)$, $K_{DT}(T)$

T, K	$K_{HD}(T)^*$, atm	$K_{HT}(T)^*$, atm	$K_{DT}(T)^*$, atm	$K_{HD}(T)^{**}$	$K_{HT}(T)^{**}$	$K_{DT}(T)^{**}$
300	18.30×10^{-22}	1.219×10^{-22}	---	3.26	2.60	3.27
400	---	---	---	3.48	2.00	3.23
500	4.639×10^{-14}	1.171×10^{-11}	---	3.62	3.26	3.04
600	---	---	---	3.77	3.45	3.01
700	---	---	---	3.76	3.40	3.07
800	---	---	---	3.83	3.68	3.01
900	---	---	---	3.91	3.76	3.01
1000	5.174×10^{-10}	2.972×10^{-10}	3.0×10^{-10}	3.93	3.91	3.00
1250	---	---	---	3.95	3.91	4.00
1500	3.109×10^{-10}	2.279×10^{-10}	2.44×10^{-10}	3.96	3.94	4.00
2000	2.611×10^{-6}	2.227×10^{-6}	2.35×10^{-10}	3.97	3.99	4.00

*Woolley, et al. (37)
**Jewes (38)

expressions for v_{HD} , v_{DT} , and v_{HT} for substitution into the first three equilibrium relations are now of the form:

$$y_{HD} = (y_H^2 y_{H_2}^{-1/2} v_{HD}^2 v_{D_2}^{-1/2} - v_{T_2}^2) / \left[\sqrt{\frac{K_{H_2}(T)}{D_{HD}^2}} + \sqrt{\frac{K_{D_2}(T)}{D_{DT}^2}} + \sqrt{\frac{K_{T_2}(T)}{D_{HT}^2}} \right]$$

To obtain a meaningful comparison with the intuitive mixing rule of Hickman, we consider a mixture of protium and deuterium only at temperatures sufficiently low and pressures sufficiently high that dissociation to form the atomic species can be neglected, but temperatures also low so that the equilibrium constant $K_{H_2}(T)$ is sufficiently small, the high-temperature limit. Based on these assumptions, solution for the molecular mole fraction yields:

$$y_{H_2} = y_H^2 / (y_H^2 + y_D^2)$$

Substitution of these results into the mixing rule yields a result identical to that of the mixing rule proposed by Hickman. To compare the mixing rule to the intuitive mixing rule of Hickman, we note that at the temperature at which Hickman's results were experimentally obtained, $K_{H_2}(T)$ has the value 0.149, and the equilibrium mixture is about 10% atomic hydrogen. The protium and deuterium isotopes are not separated to any extent, and the mixing rule is not applicable. The present results are accurate enough to allow comparison of the intuitive empirical mixing rule with the more rigorous mixing rule derived here.

ISOTOPE SWAPPING

The introduction of large amounts of protium into fusion reactor systems has been advocated to provide recovery of tritium so that lower tritium concentrations, and hence lower inventories and escape rates, can be maintained. This process, referred to as isotopic swapping, can be beneficial if:

- required capacities of tritium recovery and isotopic separation systems are not increased inordinately as a result of isotopic swapping,
- isotopic swapping significantly decreases residual tritium concentrations or enhances the effectiveness of barriers to tritium escape by mechanisms not otherwise available,
- the use of tritium recovery methods which are inherently less expensive, e.g., whose capital and/or operating costs are less, is made feasible through isotopic swapping.

These characteristics represent the basis upon which a rule of thumb for a priori assessments of potential beneficial effects of isotopic swapping on fusion reactor tritium recovery methods can be related. Many processes have been proposed for recovery and/or concentration of bred tritium from tritium-breeding blanket and coolant-loop streams, and unburned tritium from cavity exhaust streams. The effects of isotopic swapping on all these proposed processes cannot be discussed here, and we are forced to concentrate on only a few. In particular, we will discuss effects of isotopic swapping on recovery of tritium by means of tritium windows and cold trapping of liquid-metal trapping cycle and secondary coolant loops, and on escape of tritium into the biosphere by permeation through metallic barriers.

Iron-59 Swapping Effects on a Bell Metal Diffusion Barrier

Hickman¹ examined the effects of isotopic swapping on tritium-window processes for tritium extraction from liquid lithium, which possessed the following characteristics:

- Sieverts' law of solubility of hydrogen isotopes in metals and Hickman's mixing rule for dissolution of hydrogen isotopes in metals are both applicable,
- The only significant resistance to permeation through the window is the resistance to diffusion of the bulk metal of the window membrane,
- Tritium extraction is performed under steady state, isothermal conditions.

- Equilibrium with respect to isotopic interchange between diatomic molecular hydrogenic species in the gas phase is attained at all times, the high-temperature limiting values of the isotopic interchange equilibrium constants are applicable, and molecular dissociation to give the atomic species can be neglected.

The stated assumptions imply that expressions for the permeation fluxes of protium and tritium through the tritium window can be written in the forms:

$$J_{T_2} = \frac{\sigma_{T_2}(T)}{s} (D_{T_2}^{1/2} P_{T_2}^{1/2}) = \frac{\sigma_{T_2}(T)}{s} \left[y_T^0 P_{Total}^0 \right]^{1/2}$$

$$= \frac{\sigma_{T_2}(T)}{s} \left[(y_{T_2} P_{Total})^{1/2} - (y_{T_2} P_{Total}^0)^{1/2} \right]$$

$$= \frac{\sigma_{T_2}(T)}{s} \left[(y_{T_2} P_{Total})^{1/2} - (y_{T_2} P_{Total}^0)^{1/2} \right]$$

$$\sigma_{T_2}(T) = \frac{D_{T_2}(T)}{K_M} \left(\frac{2}{T} \right) \left(\frac{2}{T} \right)$$

If all these hydrogens isotopes are present and/or if dissociation of the molecular species to give the atomic species in the gas phase cannot be neglected, then the analysis will be more complex, but similar.

Now let the total pressure on the low-pressure side be negligible and consider the addition of protium to a system containing initially only tritium. Then the ratio of the tritium flux to its initial value is:

$$J_{T_2} / J_{T_2}^0 = y_T^0 (P_{Total}^1 / P_{Total}^0)^{1/2}$$

whereas the ratio of the effective system tritium inventories under the same conditions is:

$$C_T / C_T^0 = (P_{T_2}^1)^{1/2} (P_{H_2}^1 / P_{T_2}^0)^{1/2} (y_T^0 P_{Total}^1 / P_{Total}^0)^{1/2}$$

$$C_T / C_T^0 = P_{T_2}^1 / P_{T_2}^0 = y_T^0 P_{Total}^1 / P_{Total}^0$$

respectively, for a gas-phase and a liquid-metal phase constituting the principal tritium reservoir.

If the system tritium inventory is maintained constant, then, respectively, for the two cases:

$$y_T^0 P_{Total}^1 / P_{Total}^0 = y_T^0 (P_{Total}^1 / P_{Total}^0)^{1/2}$$

$$J_{T_2} / J_{T_2}^0 = y_T^0^{1/2} = 1, \quad J_{T_2} / J_{T_2}^0 = 1$$

In the first case, although the system tritium inventory has not been reduced, the tritium window area will have to be increased to remove tritium from the system at the same total rate. In the second case the isotopic sweeping is not affected the required tritium window area. Thus, depending on what type of phase represents the principal tritium reservoir, the tritium window area required for a constant total rate of tritium recovery is either increased or not affected for a constant tritium inventory, but clearly is increased if the tritium inventory is to be reduced.

If, instead, the total hydrogen isotope pressure is maintained constant, then:

$$J_{T_2} / J_{T_2}^0 = y_T^1 = \lambda$$

in both cases, while respectively:

$$C_T / C_T^0 = y_T^1, \quad C_T / C_T^0 = y_T^1^2$$

Thus, in both cases, the tritium window area required to effect the indicated inventory increase at a constant tritium recovery rate is increased by the factor $1/\lambda$, but the tritium inventory reduction is greater in the second case.

It is also clear that isotopic sweeping will either enhance or not affect the performance of barriers to tritium migration for all the cases considered. Note also that the increase in tritium window area required for removal of tritium at a constant rate when no protium is added is $1/\lambda$ when the tritium inventory level is reduced by the factor λ in either a gas phase or a liquid-metal phase.

Isotopic Swamping Effects When Surface Processes are Rate Determining

Nickman⁽³⁵⁾ and Varoni⁽³⁶⁾ examined another case of tritium recovery through a tritium window from a fluid stream that had been isotopically swamped with protium. This case involved the postulate that the permeation rate-limiting step was a slow surface process, the recombination of hydrogen isotope atoms followed by desorption from the low tritium partial-pressure side of the membrane. Adsorption-desorption processes, even if consideration is restricted to adsorption and desorption of hydrogen isotopes only, can be complex and there is a vast, often contradictory literature concerning this topic. Therefore, we can only briefly touch upon the subject. If the adsorption and dissociation step on the upstream side of the barrier is permeation rate controlling then a similar approach can be used in the analysis.

This particular step in the permeation process is carefully examined because there is evidence that, under protected thermonuclear reactor conditions, the recombination and desorption of tritium atoms may indeed be rate-limiting. For example, Gondall, et al,⁽³⁷⁾ explained the results of their experiments concerned with permeation of hydrogen isotopes from liquid lithium through niobium membranes on the basis of such a rate-limiting step.

The model for hydrogen isotope permeation which we shall consider involves the following assumptions, and approximations:

- All resistances to permeation other than those due to the desorption process are negligible.
- The partial pressures of tritium and protium on the low partial pressure side of the barrier are so small that the adsorption rate on that side of the barrier is negligible in comparison to the desorption rate.
- The simple Langmuir adsorption theory with a single class of active sites adequately describes the adsorption process.
- Dissociation in the gas phase to form the atomic hydrogen isotope species from the diatomic molecular species can be neglected.

- The adsorptive characteristics of the high and low partial-pressure barrier surfaces do not differ significantly.
- The permeation occurs under steady-state, isothermal conditions.

The essence of the elementary Langmuir theory of adsorption may be stated as follows. Adsorption of molecules from the gas phase onto solid surfaces can occur only at certain active "sites" on the surface which have an affinity for the species being adsorbed, are uniformly distributed, and have equal affinities for the adsorbed species. Each such site is assumed to be capable of adsorbing only a single atom or molecule of the adsorbed species and, because the forces that are responsible for the adsorption are short range, the adsorption at any site is assumed to be independent of whether or not an adjacent site is occupied. If the adsorption involves a dissociation, say of a diatomic molecule, such as tritium molecules, to form the atomic species, then the adsorption process is assumed to involve a reaction between the diatomic molecule and two adjacent sites on the solid surface. The number of pairs of adjacent sites, based on the above listed characteristics of the active sites, is proportional to the square of the number of unoccupied sites, which can be represented in terms of a fractional occupancy of active sites by the adsorbed species, often called the surface coverage by the adsorbed species. In terms of the surface coverage by the adsorbed species and the partial pressure of the gas phase molecular form, a kinetically third-order expression for the rate of adsorption is postulated:

$$r_{T_2\text{Ads}} = k_{T_2\text{Ads}}(T)(1-\theta_T)^2 p_{T_2}$$

Similarly, the desorption reaction is assumed to involve the kinetically second order reaction of two adsorbed atoms on adjacent active sites whose number is taken as proportional to the square of the fractional coverage and, in terms of surface coverage by the adsorbed species, a rate expression is written as:

$$r_{T_2Des} = k_{T_2Des}(T) \theta_T^2$$

Under conditions of adsorption equilibrium, the two rate expressions for adsorption and desorption can be equated and solved for θ_T to obtain:

$$\theta_T = \frac{\left[\frac{k_{T_2Ads}(T)}{k_{T_2Des}(T) P_{T_2}} \right]^{1/2}}{1 + \left[\frac{k_{T_2Ads}(T)}{k_{T_2Des}(T) P_{T_2}} \right]^{1/2}}$$

Although other, more complex theories may describe the adsorption behavior in a particular system more accurately, the results obtained with the simple Langmuir theory should be representative and may be the most accurate in some cases, e.g., see Smith.⁽³⁸⁾ Examples of other types of models for hydrogen isotope adsorption on, and desorption from metals are those of Bucur, et al.,⁽³⁹⁾ and Bucur⁽⁴⁰⁾ for adsorption of protium on palladium films. They considered multiple types of active sites and surface reactions and there are many other such treatments. The general approach to analyzing of adsorption or desorption-step-limited permeation processes by using more complex adsorption and desorption models is similar.

The listed assumptions imply that the permeation flux, when only tritium is present, is given by the Langmuir expression for the rate of desorption of tritium from the downstream barrier surface. The surface coverage in the absence of significant upstream surface adsorption and bulk barrier metal diffusive resistances is that corresponding to equilibrium with the upstream molecular tritium partial pressure. The rate expression is therefore:

$$J_{T_2} = \frac{k_{T_2Ads}(T) P_{T_2,1}}{\left\{ 1 + \left[\frac{k_{T_2Ads}(T)}{k_{T_2Des}(T) P_{T_2,1}} \right]^{1/2} \right\}^2}$$

A similar analysis based on the same assumptions when both tritium and protium are present, can be summarized as follows. The rate expressions are of the form:

$$J_{T_2} = k_{T_2Des}(T) \theta_{T_2}^2 + \frac{1}{2} k_{HTDes}(T) \theta_{T_2} \theta_{H_2}$$

and the adsorption equilibrium expressions are of the form:

$$k_{T_2Ads}(T) \theta_{T_2}^2 + \frac{1}{2} k_{HTDes}(T) \theta_{T_2} \theta_{H_2}$$

$$= k_{T_2Ads}(T) (1 - \theta_{T_2} - \theta_{H_2})^2 P_{T_2,1}$$

$$+ \frac{1}{2} k_{HTAds}(T) (1 - \theta_{T_2} - \theta_{H_2})^2 P_{HT,1}$$

The preceding adsorption equilibrium expression and the gas-phase isotopic interchange equilibrium expression can be solved for the diatomic protium and tritium and HT partial pressures in terms of the protium and tritium atom fractions in the gas phase and of the total hydrogen isotope partial pressure, although the solution may have to be obtained by numerical techniques, and the results inserted into the rate expressions. The presence of the third hydrogen isotope, deuterium, and significant dissociation of the molecular species to form the atomic species in the gas phase further complicate the problem, but the same principles are applicable.

If the tritium is isotopically swamped with protium, we have:

$$P_{H_2} \gg P_{HT} \gg P_{T_2}, \quad \theta_H \gg \theta_T$$

and the rate and adsorption equilibrium expressions become, respectively:

$$J_{T_2} = \frac{1}{2} k_{HTDes}(T) \theta_{T_2} \theta_{H_2}, \quad J_{H_2} = k_{H_2Des}(T) \theta_{H_2}^2$$

and:

$$\frac{1}{2} k_{T_2Des}(T) \theta_{T_2} \theta_{H_2} = \frac{1}{2} k_{HTDes}(T) (1 - \theta_{H_2})^2 P_{HT}$$

ARTICLE 1

Section 1. The City of... shall be a corporation of the State of... with the power to sue and be sued, to contract, to hold property, to incur and pay debts, to borrow money, to issue and sell bonds, and to do all things which natural persons may lawfully do.

Section 2. The City of... shall have and exercise all the powers and authorities of a municipal corporation, subject to the provisions of this Charter and the laws of the State of...

ARTICLE 2

Section 1. The City of... shall have and exercise all the powers and authorities of a municipal corporation, subject to the provisions of this Charter and the laws of the State of...

Section 2. The City of... shall have and exercise all the powers and authorities of a municipal corporation, subject to the provisions of this Charter and the laws of the State of...

... of liquid carbon, liquid nitrogen, and
... solid carbon-dioxide after the solvent
... and will present the results in terms
... of calculation of a single stage. The analysis
... for carbon-dioxide is given in the table in
... units.

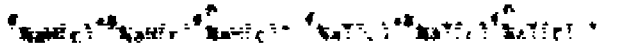
In a solid stage, the liquid must be passed
... through a series of stages, the first stage
... and the last stage being the most important
... the remaining stages. The importance of the
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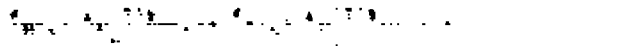
... experiments involving only a single hydrogen
... stage.

... It is assumed that the gas is to be removed
... in a single stage. The amount of gas which
... is to be removed is determined by the
... amount of gas which is to be removed.

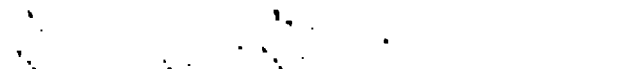


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The first part of the report discusses the general principles of the method of moments, which is a powerful tool for solving problems in structural analysis. It is particularly useful for dealing with structures that are statically indeterminate. The method involves making an assumption about the distribution of internal forces, and then using the principle of virtual work to determine the unknown forces. This process is repeated until the solution converges to the correct values.

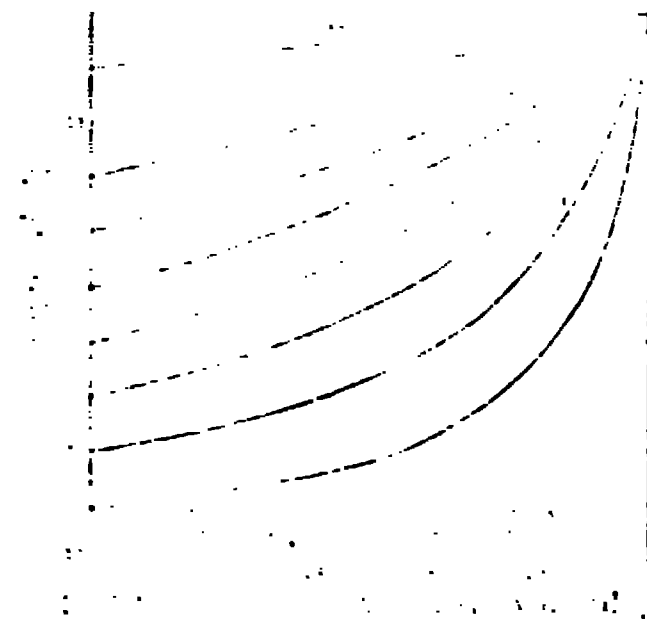
The second part of the report describes the application of the method of moments to a specific problem. The problem involves a frame structure subjected to a uniformly distributed load. The first step is to determine the fixed-end moments at the joints of the frame. This is done by assuming that the joints are fixed and calculating the moments that would be developed at these joints. The next step is to release the joints and determine the moments that would be developed at the joints if they were free to rotate. This is done by using the principle of virtual work to determine the moments that would be developed at the joints if they were free to rotate.

The final part of the report discusses the results of the analysis. The results show that the method of moments is a very effective way of solving problems in structural analysis. It is particularly useful for dealing with structures that are statically indeterminate. The method involves making an assumption about the distribution of internal forces, and then using the principle of virtual work to determine the unknown forces. This process is repeated until the solution converges to the correct values.

It is important to note that the method of moments is only applicable to structures that are statically indeterminate. For structures that are statically determinate, the method of moments is not applicable. However, for structures that are statically indeterminate, the method of moments is a very powerful tool. It allows us to determine the internal forces in a structure without having to solve a large system of equations. This is a significant advantage, especially for structures that are complex and difficult to analyze using other methods.

The method of moments is also a very intuitive method. It allows us to make an assumption about the distribution of internal forces, and then use the principle of virtual work to determine the unknown forces. This process is very straightforward and easy to understand. It is a method that can be used by anyone who has a basic understanding of structural analysis.

In conclusion, the method of moments is a very powerful tool for solving problems in structural analysis. It is particularly useful for dealing with structures that are statically indeterminate. The method involves making an assumption about the distribution of internal forces, and then using the principle of virtual work to determine the unknown forces. This process is repeated until the solution converges to the correct values.



The curves shown in the graph represent the distribution of internal forces or moments along a structure. The curves are concave up and start at a high value on the y-axis, decreasing as they move along the x-axis. The curves are labeled with different values, possibly representing different parameters or conditions. The graph is enclosed in a rectangular frame.

The curves shown in the graph represent the distribution of internal forces or moments along a structure. The curves are concave up and start at a high value on the y-axis, decreasing as they move along the x-axis. The curves are labeled with different values, possibly representing different parameters or conditions. The graph is enclosed in a rectangular frame.

- If tritium recovery system performance is determined primarily by transport rates, then isotopic swamping probably is not cost-effective and may be completely useless.
- If tritium recovery system performance is determined primarily by equilibrium considerations, then isotopic swamping may be cost effective, but there is no guarantee that this will be the case.
- Isotopic swamping will either not significantly affect or will enhance the performance of barriers to tritium escape into the biosphere for a specified tritium concentration level.

The judgment of cost effectiveness must, of course, involve consideration of changes in, e.g., required recovery system capacity, requirements for isotopic separation, benefits arising from reductions in tritium inventories and in rates of escape of tritium into the biosphere, effects on other fusion reactor systems, potential use of less expensive construction materials, or lower energy consumption, and cost of a unit of recovery system capacity.

(CONCLUSIONS)

Because of limitations imposed on the length of papers, we were unable to discuss the effects of the unavoidable presence of other hydrogen isotopes and of isotopic swamping on other fusion reactor tritium recovery systems and tritium migration mechanisms within and from fusion power plants. Examples of tritium recovery processes and migration mechanisms for which interesting analyses using similar methods can be performed include:

- the gas kinetic-limited permeation process discussed by Levin and Stickney,⁽⁵⁰⁾
- the nonequilibrium permeation process described by Shupe and Stickney,⁽⁵¹⁾
- high-temperature lithium distillation for tritium recovery suggested by Ihle and Wu,⁽³¹⁾
- the molten-salt extraction process for tritium recovery from liquid lithium described by Maroni, et al.,⁽⁵²⁾
- permeation through potential nonmetallic materials of construction for fusion reactors,

- e.g., organic polymeric materials used for, e.g., gaskets and seals, and ceramics, such as those discussed by Fowler, et al.,⁽⁵²⁾ and oxide or other corrosion layers that are potentially useful as tritium barriers, as discussed by Strehlow and Savage,⁽⁵²⁾ and
- thermal diffusion processes discussed by Pendergrass.⁽³⁰⁾

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ABBREVIATIONS

- C Concentration
- D Diffusivity
- C_0 Equilibrium concentration of the diffusing species
- C^0 Standard state free energy change
- H^0 Standard enthalpy change
- H Partition coefficient and recovery rate
- J_0 Equilibrium flux
- k Adsorption or desorption rate constant
- $k_{1,2}$ Equilibrium constants
- k_1 Solubility constant
- K_S Sieverts' law constant
- K_S^0 Sieverts' law exponential constant
- P Molecular weight

N_{Na}

N_{Na}^0

N_{Na}

P.P

K_S

r

R

S^0

T

r^0

X

V, V'

Greek Letters

- δ Trifluoromethane flow rates
- Pressure
- Sieverts' law exponential constant
- Absorption or desorption rate
- Ideal gas constant
- Standard entropy change
- Absolute temperature
- Trifluoromethane recovery rate
- Mole fraction
- Mole, atom fractions
- Trifluoromethane window or permeation barrier thickness
- Activity coefficient
- Fractional surface coverage by adsorbed species
- Trifluoromethane concentration level
- Density
- Permeability